

[Volume 18](https://jmstt.ntou.edu.tw/journal/vol18) | [Issue 6](https://jmstt.ntou.edu.tw/journal/vol18/iss6) Article 9

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Recommended Citation

Wang, Helsin; Chang, Ta-Peng; and Wang, Jin-Jun (2010) "RESPONSE ANALYSIS OF CONCRETE PILES SUBJECTED TO LATERAL IMPACT," Journal of Marine Science and Technology: Vol. 18: Iss. 6, Article 9. DOI: 10.51400/2709-6998.1943

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RESPONSE ANALYSIS OF CONCRETE PILES SUBJECTED TO LATERAL IMPACT

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Key words: dispersive flexural waves, three-dimensional wave theory, resonant frequency, phase velocity.

ABSTRACT

One-dimensional wave theory is assumed to interpret longitudinal waves impulse response or impact echo testing results in concrete piles. Transient flexural waves tests are alternatively introduced to evaluate piles where the tops are inaccessible for conventional longitudinal wave testing. The resulting interpretation for flexural waves based on one-dimensional wave theory is improper due to various velocities at low frequencies. In this research, treating a pile as a waveguide, three-dimensional wave theory provides the dispersion relation between phase flexural velocity and frequency. Using the resonance solutions and first-mode flexural wave curve is a simple and efficient method to find the resonant frequency and phase velocity by knowing pile dimension and flexural IR results. The phase velocities at resonant frequencies measured in several flexural wave tests have good matches with the first-mode wave theoretical predictions at frequencies below 3,000 Hz. The boundary conditions for practical feasibility are discussed for different *in-situ* conditions.

I. INTRODUCTION

Non-destructive testing (NDT) techniques have been applied to construction quality control for drilled shafts and concrete piles for more than forty years. Both the impulse response (IR) and impact echo (IE) methods are two most common techniques used to non-destructively evaluate the integrity of concrete piles in field. The in-service conditions of existing deep foundations, e.g., the evaluations of dimensions and continuity of pile foundations, usually have to been qualitatively evaluated by these two most popular methods [1, 3, 6, 12, 14].

Transient-state longitudinal (compressive) waves are commonly induced in conventional IR or IE tests at the top of pile

foundations by a hammer while the top of piles is accessible to install testing equipment. When the top of concrete piles is inaccessible for testing, an approach of transient-state flexural wave induced by the lateral impact is an alternative, particularly applicable in the partially exposed piles of overwater structures, such as wharf and bridge [5, 8, 9, 11, 18-21].

Normally the constant (non-dispersive) wave velocity onedimensional (1-D) wave theory is applied to interpret results of longitudinal wave IR or IE tests in concrete piles [4, 10, 21]. Using spectrum information can efficiently evaluate pile lengths or defect positions. On the contrary, the flexural wave velocities are varied with frequencies at low frequencies, overlapping with the induced frequency range in the IR or IE tests [5, 18, 23]. As a result, the feasibility of flexural wave with IR or IE tests is currently limited at high frequencies, especially improper for concrete piles with a high material attenuation [13, 21, 23].

Using the 1-D wave theory may be insufficient to provide proper interpretation of transient flexural wave experimental results at low frequency. To efficiently apply the responses at the most excited frequency range in the flexural IR tests, the dispersion relation between phase velocity and frequency is introduced by using the three-dimensional (3-D) wave theory while the concept of a waveguide is applied to concrete piles. The phase velocities at resonant frequencies in flexural wave IR or IE tests compare with the 3-D flexural wave theoretical predictions. The practical application and limitation of phase velocities at resonant frequencies are also discussed in detail.

II. ONE-DIMENSIONAL WAVE THEORY

Consider that a given slender pile is subjected to the IR or IE testing, the induced longitudinal or flexural wavelength is normally much longer than the pile dimension, usually pile diameter. If the velocity of the induced stress waves is regarded as a constant, not a function of frequency, and the crosssectional displacement is uniform, the 1-D wave theory is intuitively assumed to interpret these testing results [10, 21, 22].

Figure 1 schematically shows the typical set for a traditional IR or IE test, called as the longitudinal IR or IE test herein, that includes a modal hammer, which produces an impact at the pile top or cap, and one or several fixed geophones or accelerometers at the pile top, which measure the surface vibration reflected from either the bottom of the pile or

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Fig. 1. Schematic arrangement of longitudinal impulse response testing.

defects in the pile along a longitudinal path parallel to the axial direction [1, 6, 12, 14, 17]. Both the force and velocity or acceleration signals are measured and digitally recorded in a portable data acquisition system. These time-domain signals are converted to the frequency domain or spectrum with a fast Fourier transfer (FFT) for further analysis. Conventionally, the bar wave velocity, C_{bar} , in a concrete pile can be calculated from either the time domain or frequency domain data. In the time domain, the bar wave velocity, C_{bar} , is computed by equation

$$
C_{bar} = \frac{2\ell}{\Delta t_{avg}} \tag{1}
$$

where ℓ is the distance between impact point and reflecting sources, either the bottom of a pile or a defect, and ∆*tavg* the average traveling time between the impact point and reflection sources [1, 3, 4, 12]. The computed wave velocity based on the time domain information usually represents a group velocity, namely the velocity of energy transportation of a wave [4]. In the frequency domain, the bar wave velocity, C_{bar} , is computed by

$$
C_{bar} = 2\ell \times f_{peak} \tag{2}
$$

where *fpeak* is the maximum peak frequency [3, 4, 14]. The common bar wave velocity, C_{bar} , of a sound concrete pile has a wide range from 3,500 to 4,500 m/sec [3, 12]. If knowing the bar wave velocity, C_{bar} , and second dominant peak frequency, *fdefect*, one can sketchily identify the significant defect position, ℓ_{defect} , by re-grouping (2) as the following:

$$
\ell_{defect} = \frac{C_{bar}}{2f_{defect}}
$$
 (3)

The experimental procedures of the flexural wave method are similar to those of the conventional longitudinal IR or IE method. Figures 2(a) and 2(b) show two possible flexural IR or

(a) Pile head exposed to the open space

(b) Pile head connected to a beam

Fig. 2. Two schematic arrangements of flexural impulse response or impact echo testing.

IE tests conducted on a partially exposed concrete pile [5, 7-9, 16-20]. The former needs to remove a portion of topsoil surrounding the pile, while the latter is performed for the concrete shafts or piles connected to the docks, bridges, or wharves with lack of access to their tops. The transient-state flexural waves, traveling downward and upward along the pile, are induced by lateral impact on the free side of the concrete pile and recorded by two sensors mounted to either the opposite sides on the periphery or the same sides of the pile. The data of acceleration signals at time domain are stored in a portable data acquisition and converted to the frequency domain with FFT.

Usually, the apparent flexural wave velocity, C_s , in a concrete pile is computed in the time domain by

$$
C_s = \frac{D}{\Delta t} \tag{4}
$$

where *D* is the distance between two sensors and ∆*t* is the traveling time between two sensors [5, 7-9, 16-20]. The flexural wave velocity, C_s , in a sound concrete pile commonly

Fig. 3. Phase velocity curves of L(0,1) and F(1,1) in a 1-m diameter pile with shear wave velocity of 2,500 m/sec.

varies from 2,100 to 2,700 m/sec [3], which can be applied to identify the pile lengths [7-9, 11, 16-20, 22] and localize the flaw positions [5, 18].

Frequency analysis currently is not recommended for the flexural IR or IE testing due to its dispersive characteristics, i.e., the wave velocity varied with frequency [5, 7-9, 18, 21]. A numerical analysis based on the Timoshenko beam theory suggested that 1-D wave theory is acceptable in the flexural IR or IE testing for the resonant frequency modes higher than eight such that the wave velocity is approaching a constant value [23]. The analytic and experimental researches based on the three-dimensional wave theory supported the findings of non-dispersive flexural waves at higher frequencies in an embedded concrete pile with the same manner [16, 21].

III. THREE-DIMENSIONAL WAVE THEORY

1. Longitudinal Waves

The stress waves traveling in piles, called three-dimensional (3-D, guided) waves, include different clusters of waves arising from incidences and reflections of compression, shear, and surface waves along the boundaries of piles. The real wave velocity for such case is a function of frequency, and the displacement magnitudes vary along the pile cross-section. Theoretical and experimental investigations of the longitudinal and flexural waves in the prototype concrete piles embedded in a sand fill have been derived and verified, respectively [4, 13, 16, 21].

The maximum frequency generated generally reaches as high as 2 to 6 kHz by hammer impacts used in longitudinal IR tests [10, 12]. Its frequency range is located on the wave branch $L(0,1)$, the first mode in the family of longitudinal

modes with no radial displacement, derived from the guided wave theory, which is dispersive and varied with frequency as shown in Fig. 3. The phase velocity remains approximately a constant value around the bar wave velocity for frequencies below 2,000 Hz, but approaches the Rayleigh wave velocity at frequency beyond 5,000 Hz. Thus, a transient phase velocity monotonically decreasing from bar wave velocity to the Rayleigh wave velocity develops at frequency between 2,000 to 5,000 Hz.

A transient-state (impact) wave can be considered as a superposition of many steady-state (3-D) waves [4]. The predominant waves excited by IR methods correspond to the low frequency range of the lowest branch of longitudinal guided waves in concrete piles and drilled shafts [4, 10]. Due to the limitation of excited longitudinal waves within a low and narrow frequency range associated with IR tests, the wave velocities do not change much with frequency in concrete piles. Its longitudinal wave velocity in these narrow frequency ranges usually can be regarded as a constant that does not change with frequency [10] so as to the phase velocity in such a frequency range. Once the assumption of constant longitudinal wave velocity is accepted, the concept of 1-D waves, and not dispersive waves, is applied in general. Therefore, the 1-D wave theory used by the IR techniques is one special case of the 3-D wave theory at frequencies below 500 Hz and 6,000 Hz for pile diameters 2.4 m and 0.3 m, respectively [21].

2. Flexural Waves

When the flexural waves on piles are introduced by a modal hammer, the most excited flexural waves are located within a low frequency range as those with a longitudinal IR method. Many experimental and theoretical investigations demonstrated the existence of dispersion of flexural waves in timber or concrete piles [5, 11, 17-23]. The phase velocity of flexural waves is no longer regarded as a constant value but various at low frequencies. The analytic solution of 3-D flexural wave theory suggests that the assumption of constant flexural wave velocity is acceptable only at frequencies of above 1,500 Hz to 10 kHz for pile diameter of 2.4 to 0.3 m, respectively [21]. This fact emphasizes the need to measure the flexural wave velocity directly at low frequencies, since it is not reasonable to assume a constant propagation velocity, as is commonly done for the longitudinal wave tests [11]. Using the 1-D wave theory may be insufficient to provide proper interpretation of experimental results [15].

Treating a pile as a wave guide, the lowest flexural wave branch $F(1,1)$, the first mode in the series of flexural modes with one displacement cycle in the radial direction of the waveguide, is located within the most possible induced frequency range in the flexural wave tests [11]. Figure 3 also shows a typical dispersive phase velocity curve of branch $F(1,1)$ for flexural waves traveling in a 1-meter diameter pile with a shear wave velocity of 2,500 m/sec. At low frequencies, the velocity value increases monotonically from zero to the Rayleigh wave velocity, 2,260 m/sec with frequency. For

Fig. 4. Phase velocities of longitudinal waves in a drilled shaft [4].

frequency beyond 2,500 Hz, the flexural wave velocity approaches the Rayleigh wave velocity. For such frequency range, the flexural wave velocity becomes apparently constant and non-dispersive with frequency, instead. The existence of plateaus for branch $F(1,1)$ represents that the assumption of a constant velocity for flexural waves is valid only at high frequencies. On the other hand, the wave velocity of flexural waves induced by IR methods is dispersive at low frequencies by using 3-D wave theory. The constant velocity 1-D wave theory cannot be used in explaining flexural wave results at low frequencies.

3. Connection between 3-D Waves and Longitudinal and Flexural Wave IR Tests

When conducting IR tests, either longitudinal waves or flexural waves, in concrete piles or drilled shafts, the incident waves on the pile top are supposed to be reflected on the pile bottom. The measured time is long enough to detect multiple reflections from the bottom, and arises resonance responses in concrete piles or shafts.

The resonant solutions of longitudinal IR tests in a finite embedded pile have been developed based on 1-D wave theory [4]. The longitudinal wave velocity in a low and narrow frequency ranges usually can be regarded as a constant that does not change with frequency. It is a special situation to find a regular resonant frequency from the longitudinal wave spectrum. It is the fundamental to use the resonant solutions and subsequently find its resonant frequency and phase velocity.

When the pile length is known, the phase velocity of longitudinal or flexural waves can be derived based on such a method. The experimental phase velocities are superimposed on 3-D wave theoretically-determined longitudinal wavecurve $L(0,1)$. The IR experimental results are expected to lie along

these theoretical prediction curves. This resonance technique has been already successfully applied to the longitudinal wave experimental results in concrete piles in few years ago. The computed phase velocity is given in Fig. 4 for the longitudinal wave results with an assumption of displacement CB conditions [4]. The results show a good trend between the 3-D wave theoretical curve and IR experimental results. However, the same resonance technique has not yet been used to predict the resonant frequency and phase velocity in flexural wave experimental results due to changing velocity with frequency.

For flexural tests, the relevant numerical researches recommend that the frequency difference between two consecutive resonant frequencies at high frequencies approaches a constant value in piles, and the resonant frequencies measured from the flexural IR tests can be applied to evaluate pile lengths [19, 20, 23]. To the contrary, the dispersive flexural velocity arises an irregular frequency difference between two consecutive resonant frequencies at low frequencies which are the most excited frequency range in the flexural IR tests. This indicates that the resonant solutions in a finite embedded pile cannot be used in the same method as the situation in longitudinal wave experiments at low frequencies.

There are several disadvantages to use numerical analysis for flexural IR tests in piles. Relatively high attenuation in geomaterials, especially the embedded concrete piles herein, at high frequencies undermines the advantage of applying constant resonance difference [16, 21], not as expected in numerical analysis. The response at high frequencies is not easily excited and measured in concrete piles with flexural IR tests. In addition, most of these researches use numerical finite element method to circumvent the drawback in order to find the resonant frequency and phase velocity. A number of parameters about piles and site situation, including pile dimension, mechanical property, pile-to-soil interaction model, and engineering properties of soils, should be provided and determined in advance.

The whole process of data collection and analysis is timeconsuming and complicated by using numerical analysis. In this research, the resonance technique and first-mode flexural wave curve $F(1,1)$ are used to find the resonant frequency and phase velocity based on known pile dimension and flexural wave testing only. The experimental phase velocities are superimposed on 3-D wave theoretically-determined curve F(1,1) for flexural waves.

To efficiently apply the resonance responses at the most excited frequency range in the flexural IR tests, the resonance solutions of stress waves propagating in a finite embedded pile with a length of *L* can be found as [2]

$$
\lambda_d = \frac{2L}{n} \tag{5}
$$

$$
\lambda_s = \frac{4L}{2n - 1} \tag{6}
$$

$\frac{1}{2}$						
Property	Pile A	Pile B	Pile C			
Diameter	0.355 m	0.254 m	0.308 m			
Length, L	2.43 m	2.22 m	2.4 m			
Bar wave velocity, C_{bar}	3900 m/sec	4150 m/sec	4097 m/sec			
Shear wave velocity, C_s	2530 m/sec	2625 m/sec	2592 m/sec			
Poisson's ratio, ν	0.2	0.25	0.28			

Table 1. Physical properties of prototype concrete piles [4, 16].

Fig. 5. Displacement responses for a pile subjected to a lateral impact.

respectively, where λ_d is the resonant wavelength of waves with a displacement control boundary (CB), λ_s the resonant wavelength of waves with a stress CB, and *n* the positive integral resonant number. Each resonant frequency, f_n , is corresponding to a resonant number, *n*, and a resonant wavelength.

The phase velocity of longitudinal or flexural waves in a finite length pile can be measured in the frequency domain by means of an IR method. Generally, several resonant frequencies can be observed in IR testing results in concrete piles. The phase velocity, C_{ph} , for each resonant frequency, f_n , is defined as

$$
C_{ph} = \frac{\omega}{k} = \lambda \times f_n \tag{7}
$$

where ω is the angular frequency equal to $2\pi f_n$, *k* is the wave number equal to $2\pi/\lambda$, and λ is the resonant wavelength, either λ_d or λ_s . The computed wave velocity based on the frequency domain information usually represents a phase velocity, i.e., the velocity of a particle on a wave [4, 11, 22].

4. Vibration Responses of Flexural IR Tests

Impacts on the side of a pile are suggested as a method to induce transient-state flexural waves [4, 16-20]. The displacement responses for a fully or partially embedded pile subjected to a lateral impact at pile top are schematically shown in Fig. 5. The orthogonal cylindrical coordinate, $r-\theta-z$, system and the impact position (assumed 0°) are given in the left side of Fig. 5. The expected deformation profiles from the side and top views are also shown. A warping-type displacement pattern occurs at top of the pile.

The displacement response at point A ($\theta = 0^{\circ}$), the impact position, undergoes negative radial displacement, *ur*, and negative axial displacement, u_z , with no angular displacement. The displacement response at point A' ($\theta = 180^{\circ}$) undergoes positive radial and positive axial displacements with no angular displacement, u_{θ} . In practice, the displacements in the angular/tangential direction are normally negligible and the measurable displacements at positions 0° and 180° occur in the radial and longitudinal directions.

The displacements at points B ($\theta = 90^{\circ}$) and B' ($\theta = 270^{\circ}$) undergo no radial and axial displacements, but positive and negative angular displacements, respectively. In practice, the displacements in the radial and longitudinal directions are normally negligible. At positions 90° and 270°, the maximum displacement occurs in the angular/tangential direction. The analytic displacement responses for the first branch $F(1,1)$ of the 3-D (steady-state) waves represent the displacement responses of transient flexural waves [20, 23].

IV. FLEXURAL WAVE IR TESTS ON CONCRETE PILES

1. Prototype Concrete Piles

In this study, the flexural and longitudinal wave IR tests are conducted to three prototype concrete piles with different diameters of 0.254 to 0.355 m embedded in sandy fill (as the left side of Fig. 2). The physical properties of the three concrete piles are listed in Table 1 [4, 16].

The bar wave velocity and shear wave velocity are computed from the time-domain vibration responses. The Poisson's ratio in concrete is computed by the following equation:

$$
V = \frac{C_{bar}^2}{2C_s^2} - 1\tag{8}
$$

The lateral impacts were conduced at point A (assumed 0°) with a modal hammer as that in Fig. 5. The vibration responses were measured with two triaxial accelerometers on the pile side after moving part of surrounding soils from the ground. The location and orientation of the transducer is consistent with the deformation induced by the impact. A typical FFT spectrum in the angular/tangential direction at positions 90° and 270° for pile A is shown in Fig. 6.

A set of resonant peaks is found at frequencies below 3,000

Table 2. Phase velocity computation using measured resonant frequency in prototype concrete pile A at measurement positions 90° and 270° (angular/tangential direction) with stress and displacement control boundaries.

	Position at 90° Position at 270°		Stress control boundary		Displacement control boundary	
			C_{ph}		\mathbf{C}_{ph}	
n	Jn	J _n	90°	270°	90°	270°
	51	51	496	496	248	248
◠	224	255	724	826	543	619
3	549	529	1067	1029	889	858
4	990	990	1375	1375	1203	1203
	1365	1373	1474	1482	1326	1334
6	1839	1851	1625	1636	1490	1499
	2314	2333	1730	1745	1606	1620
\circ δ	2706	2765	1753	1792	1644	1680

Note 1: *n* is the resonant number, f_n is the resonant frequency, and C_{ph} is the phase velocity.

Note 2: Units of resonant frequency and phase velocity are Hz and m/sec, respectively.

Note 1: *n* is the resonant number, f_n is the resonant frequency, C_{ph} is the phase velocity, and N/I indicates not identified. Note 2: Units of resonant frequency and phase velocity are Hz and m/sec, respectively.

Fig. 6. Flexural impact results on a prototype concrete pile (Solid and dashed lines represent measurement positions at 90° and 270°, respectively).

Hz. The first eight measured resonant frequencies are indicated in Table 2. The resonant number, *n*, is listed in the first column of Table 2. The identified resonant frequencies are listed in the second and third columns of Table 2 at positions 90° and 270°, respectively. Similar resonant frequency values are observed by two accelerometers. In view of resonance principle, since the resonant frequencies for a wave traveling in a concrete pile are independent on their measurement positions, the same resonant frequencies at different measurement positions are theoretically expected. The phase velocities corresponding to resonant frequency are computed in the rest columns of Table 2 by using (7) for stress and displacement CBs. The similar phase velocity computations using resonant frequencies measured at positions 0° and 180° in the radial and longitudinal directions are also listed in Tables 3 and 4, respectively.

The flexural wave IR experimental results are superimposed on the theoretical prediction curve F(1,1) illustrated in Figs. 7 to 9. Since the induced frequency generated by a

	Position at 0°	Position at 180°	Stress control boundary		Displacement control boundary	
			C_{ph}		C_{ph}	
n	Jn	J_n	0°	180°	0°	180°
	N/I	N/I	N/I	N/I	N/I	N/I
	220	210	713	680	535	510
3	515	530	1001	1030	834	858.6
	915	915	1271	1271	1112	1112
	1365	1380	1474	1490	1327	1341
6	1870	N/I	1652	N/I	1515	N/I
	2285	N/I	1708	N/I	1586	N/I
8	2714	2750	1759	1782	1649	1671

Table 4. Phase velocity computation using measured resonant frequency in prototype concrete pile A at measurement positions 0° and 180° (longitudinal direction) with stress and displacement control boundaries.

Note 1: *n* is the resonant number, f_n is the resonant frequency, C_{ph} is the phase velocity, and N/I indicates not identified. Note 2: Units of resonant frequency and phase velocity are Hz and m/sec, respectively.

Fig. 7. Phase velocities of flexural waves a prototype concrete pile at measurement positions 90° and 270° (angular/tangential direction).

Fig. 8. Phase velocities of flexural waves a prototype concrete pile at measurement positions 0° and 180° (radial direction).

modal hammer ranges less than 4,000 Hz on concrete piles, the resulting phase velocities can be described by flexural wave curve $F(1,1)$. The phase velocities at identified resonant frequencies monotonically increase with frequency and have similar velocity trends with 3-D wave theoretical predictions. The analysis results agree with the conclusions published in [4, 11, 16]. Most of the computed phase velocities with a stress CB lie on along the theoretical prediction curve. The computed phase velocities with a displacement CB follow the shape of the theoretical curve. This indicates that phase velocities with a stress CB have a better curve matching than those with a displacement CB. For inducing a uniform lateral deformation is much more difficult than inducing a uniform shear stress at a pile top by a hammer, the mechanical condition set as a stress CB at the pile top is more reasonable for lateral impacts on embedded concrete piles.

2. Precast Concrete Piles

A cast-in-place concrete dock is supported by a transfer beam and pile system. Each 254-mm-square precast concrete pile is 15.25 m in length with design compressive strength of 27.6 MPa. Flexural and longitudinal wave IR tests were conducted on 5 in-service precast concrete piles [8, 9]. The experimental arrangement of flexural waves induced on these 5 concrete piles is shown as the right side of Fig. 2. Table 5 describes the lengths of piles and material properties, bar wave velocity, shear wave velocity, and Poisson's ratio of concrete.

The orientation of two sensors is consistent with the deformation induced by lateral impacts and the displacement responses along the pile side are like the responses at point A (assumed 0°) in Fig. 5. The maximum displacement theoretically occurs in the radial (impact) direction. Their displacements in the tangential and longitudinal directions can

Fig. 9. Phase velocities of flexural waves a prototype concrete pile at measurement positions 0° and 180° (longitudinal direction).

Table 5. Summary of propagation velocities for 5 precast concrete piles [8, 9].

Pile	Length (m)	C_{bar} (m/sec)	C_s (m/sec)	
	15.76	4026	2623	0.18
2	16.0	4300	2806	0.18
3	15.99	4026	2593	0.19
	15.76	4026	2593	0.21
	16.0	4184	2562	0.20

Note 1: Lengths of piles indicate the total length of the flexural waves propagating through a precast pile, transfer beam, and castin-place concrete dock.

be negligible due to their superstructure constraints. A typical radial-direction acceleration spectrum obtained from a flexural wave IR test in a dock concrete pile is shown in Fig. 10.

The resonant frequencies corresponding to the first 39 resonant numbers are indicated in Table 6. The resonant number is listed in the first column of Table 6. The similar identified resonant frequencies on upper and lower sensors in the radial direction are listed in the second and third columns of Table 6, respectively. Some resonant frequencies are not identified in the FFT spectrum. The possible causes are an FFT spectrum with a lower resolution due to a lower sampling rate in testing or the FFT responses masked by adjacent peaks. The similar phase velocities with stress and displacement CBs are computed in the rest columns of Table 6.

A typical relation chart for phase velocity and frequency measured by flexural wave IR testing is plotted in Fig. 11. The computed phase velocities at resonant frequencies monotoni-

Fig. 10. Transient-state flexural wave response spectrum in a partially exposed concrete pile.

Fig. 11. Phase velocities of flexural waves in a partially exposed concrete pile measured by two sensors.

cally increase with frequency for upper and lower sensors. There is a good match for phase velocities between 3-D wave theoretical prediction curves and flexural wave IR testing results for frequency up to 3,000 Hz. The analysis results agree with the conclusions published in [4, 11, 16]. Both computed phase velocity points for stress and displacement CBs fall around the theoretical prediction curve (see Fig. 12). The result implies that the real boundary condition lies in between stress and displacement controls for dock concrete piles.

3. Practical Implication

The analysis results in embedded prototype concrete piles

Note 2: C_{bar} is the bar wave velocity, C_s is the shear wave velocity, and v is the Poisson's ratio of concrete.

			Stress control boundary C_{ph}		Displacement control boundary	
	Upper sensor	Lower sensor			\mathcal{C}_{ph}	
\boldsymbol{n}	\mathfrak{f}_n	\mathfrak{f}_n	Upper sensor	Lower sensor	Upper sensor	Lower sensor
$\mathbf{1}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$N\!/\!I$	$\rm N/I$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\rm N/I$
$\overline{2}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	30	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	640	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	480
3	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$N\!/\!I$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$
$\overline{4}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$80\,$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	731	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	640
5	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	160	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1138	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1024
6	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	220	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1280	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1173
$\boldsymbol{7}$	260	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1280	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1189	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$
$\,8\,$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	300	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1280	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1200
9	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	410	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1544	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1458
10	450	$N\!/\!I$	1516	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1440	$N\!/\!I$
11	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	540	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1646	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1571
12	630	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1753	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1680	\mathbf{N}/\mathbf{I}
13	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	690	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1766	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1698
14	790	780	1873	1849	1806	1783
15	890	890	1964	1964	1899	1899
16	1000	1000	2064	2064	2000	2000
17	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1110	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2153	$N\!/\!I$	2089
18	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$N\!/\!I$
19	1230	1240	2128	2145	2072	2088
20	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$
$21\,$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1440	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2248	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2194
22	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$N\!/\!I$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$N\!/\!I$
23	1540	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2190	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2143	$N\!/\!I$
24	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1670	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2274	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2227
25	1760	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2299	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2253	$N\!/\!I$
26	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	1850	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2322	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2277
$27\,$	1940	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2343	$N\!/\!I$	2299	$N\!/\!I$
28	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$N\!/\!I$	$N\!/\!I$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$N\!/\!I$
29	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2150	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2414	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2372
30	2180	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2365	$N\!/\!I$	2325	$N\!/\!I$
31	\mathbf{N}/\mathbf{I}	2300	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	2413	\mathbf{N}/\mathbf{I}	2374
32	2410	2410	2448	2448	2410	2410
33	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$	\mathbf{N}/\mathbf{I}	$\ensuremath{\text{N}}\xspace/\ensuremath{\text{I}}$
34	2560	2540	2445	2426	2409	2391
35	2670	2680	2476	2486	2441	2450
$36\,$	N/I	N/I	N/I	N/I	N/I	N/I
37	N/I	N/I	N/I	N/I	N/I	N/I
38	2900	2900	2475	2475	2442	2442
39	2980	N/I	2477	N/I	2445	N/I

Table 6. Phase velocity computation using measured resonant frequency in the radial direction in Pile 2.

Note 1: *n* is the resonant number, f_n is the resonant frequency, C_{ph} is the phase velocity, and N/I indicates not identified. Note 2: Units of resonant frequency and phase velocity are Hz and m/sec, respectively.

and precast concrete piles provide a 3-D flexural wave based interpretation in explaining irregular frequency difference between two consecutive resonant frequencies at low frequencies [23]. Knowing the designate pile length, *L*, and apparent flexural wave velocity, *Cs*, the flexural wave IR or IE tests can be applied to evaluate *in-situ* pile lengths or defect positions if

the computed phase velocities corresponding to measured resonant frequencies cannot completely superimpose on 3-D flexural wave prediction curves.

In addition, the identification degree of resonant frequency is affected by length-to-diameter ratio for the higher lengthto-diameter ratio is and the lower identification degree of

Fig. 12. Average phase velocities of flexural waves in a partially exposed concrete pile.

resonant frequency is. The identification degree of resonant frequency for embedded prototype concrete piles with lengthto-diameter ratios of 6.8 to 8.7 is higher than that for dock precast concrete piles with length-to-diameter ratios of around 17.

V. CONCULSIONS

When the top of piles is conveniently accessible to install testing equipment, longitudinal waves are commonly induced to perform IR or IE tests in concrete piles. However, when the pile top is inaccessible as a result of intervening structure, for example transfer beams, or superstructures, only a small portion of the longitudinal wave energy will actually reach the concrete pile. A flexural wave IR or IE test will be a more efficient NDT technique to conduce integrity evaluation in piles. Transient flexural waves are induced by lateral impact on a partially exposed concrete pile side. Based on the analysis results, the following conclusions can be drawn:

- 1. The resonance technique and first-mode flexural wave curve $F(1,1)$ are used to find the resonant frequency and phase velocity by knowing the pile dimension and flexural IR testing only. It is a simpler and more efficient method than that by using numerical finite element analysis.
- 2. The computed phase velocities at measured resonant frequencies have a good agreement with 3-D flexural wave theoretical prediction curves and monotonically increase with frequency at frequencies below around 3,000 Hz.
- 3. The boundary condition also affects the computed phase wave velocity. For flexural wave IR or IE tests conduced in embedded concrete piles, a stress control boundary can properly represent the mechanical condition for its phase velocities have a better matching with 3-D flexural wave

theoretical prediction curves than those with a displacement CB. The real boundary condition lies in between stress and displacement controls for partially exposed concrete piles supporting a dock.

- 4. In practical, if the designate length and apparent flexural wave velocity are known in a pile, using flexural wave IR or IE tests can evaluate its *in-situ* dimension characteristics. When the computed phase velocities of resonant frequencies cannot completely superimpose on 3-D flexural wave theoretical curves, the occurrence of breakage and/or major defects could be expected in a pile.
- 5. For the identification degree of resonant frequency in a flexural wave spectrum, piles with a lower length-to-diameter ratio have a higher value than that for those with a higher length-to-diameter ratio.

APPENDIX

Stress waves propagate along a finite pile with a length of *L* fully or partially embedded in uniform soils/rocks. Assuming the position, *z*, equal to zero at a pile top, the pile is represented in an orthogonal cylindrical coordinate *r*–θ–*z* system. The displacement of 1-D longitudinal and flexural wave propagation in a pile at any position *z* can be expressed as

$$
u_z(z,t) = A_1 \exp[i(\omega t - kz)] + A_2 \exp[i(\omega t + kz)] \qquad (A-1)
$$

$$
u_r(z,t) = A_3 \exp[i(\omega t - kz)] + A_4 \exp[i(\omega t + kz)] \qquad (A-2)
$$

respectively, where ω is the angular frequency equal to $2\pi f$, *t* is the time, *k* is the wave number equal to $2\pi/\lambda$, A_1 and A_2 are the displacement amplitude of longitudinal waves, A_3 and A_4 are the displacement amplitude of flexural waves, *f* is the frequency, and λ is the wavelength.

In general, the pile top is allowed to take a relatively higher deformation as a pile is subjected to impacts at its top. The mechanical condition on a pile bottom is often regarded as having a very low deformation. Such boundary conditions are usually simulated as one-free-end and one-fixed-end boundary conditions in a pile. Usually, the displacement is assumed as zero at the fixed-end boundary condition. The wave modes in a finite length pile are reflected without phase change as a wave of the same type but traveling in the opposite direction [4]. However, there are two possibilities for the free-end boundary condition, a given displacement or given stress boundary condition. In this research, such two boundary conditions are simplified as displacement control boundary (CB) and stress CB, respectively.

For longitudinal waves, the displacement CB at a pile top is

$$
u_z(0,t) = A_1 \exp(i\omega t) + A_2 \exp(i\omega t) = A_{0z} \exp(i\omega t)
$$
 (A-3)

where A_{0z} is the displacement amplitude of incident longitu-

dinal waves. The displacement CB at a pile bottom is

$$
u_z(L, t) = A_1 \exp[i(\omega t - kL)] + A_2 \exp[i(\omega t + kL)] = 0 \quad (A-4)
$$

The displacement solution of longitudinal waves is

$$
u_z(z,t) = \frac{A_{0z} \sin[k(L-z)]}{\sin(kL)} \exp(i\omega t)
$$
 (A-5)

This indicates that the amplitude of the displacement is infinite when its denominator is equal to zero [2]. The evaluation of resonance solutions for longitudinal waves is summarized as

$$
kL = n\pi \tag{A-6}
$$

where *n* is the positive integral resonant number.

For longitudinal waves, the stress CB at a pile top is

$$
\sigma_{zz}(0,t) = T_{0z} \exp(i\omega t) = E \frac{\partial u_z}{\partial z}
$$
 (A-7)

where T_{0z} is the stress amplitude of incident longitudinal waves. The stress CB at a pile bottom is

$$
u_z(L,t) = A_1 \exp[i(\omega t - kL)] + A_2 \exp[i(\omega t + kL)] = 0
$$
 (A-8)

The displacement solution of longitudinal waves is

$$
u_z(z,t) = -\frac{T_{0z} \sin[k(L-z)]}{kE \cos(kL)} \exp(i\omega t)
$$
 (A-9)

This indicates that the amplitude of the displacement is infinite when its denominator is equal to zero [2]. The evaluation of resonance solutions for longitudinal waves is

$$
kL = (2n-1)\pi \tag{A-10}
$$

For flexural waves, the displacement CB at a pile top is

$$
u_r(0,t) = A_3 \exp(i\omega t) + A_4 \exp(i\omega t) = A_0 \exp(i\omega t) \quad (A-11)
$$

where A_{0r} is the displacement amplitude of incident flexural waves. The displacement CB at a pile bottom is

$$
u_r(L,t) = A_3 \exp[i(\omega t - kL)] + A_4 \exp[i(\omega t + kL)] = 0 \quad (A-12)
$$

The displacement solution of flexural waves is

$$
u_r(z,t) = \frac{A_{0r} \sin[k(L-z)]}{\sin(kL)} \exp(i\omega t)
$$
 (A-13)

This indicates that the amplitude of the displacement is infinite when its denominator is equal to zero [2]. The evaluation of resonance solutions for flexural waves is expressed as

$$
kL = n\pi \tag{A-14}
$$

For flexural waves, the stress CB at a pile top is

$$
\sigma_{rr}(0,t) = T_{0r} \exp(i\omega t) = G \frac{\partial u_r}{\partial z}
$$
 (A-15)

where T_{0r} is the stress amplitude of incident flexural waves. The stress CB at a pile bottom is

$$
u_r(L,t) = A_3 \exp[i(\omega t - kL)] + A_4 \exp[i(\omega t + kL)] = 0
$$
 (A-16)

The displacement solution of flexural waves is

$$
u_r(z,t) = -\frac{T_{0r} \sin[k(L-z)]}{kG \cos(kL)} \exp(i\omega t) \tag{A-17}
$$

This indicates that the amplitude of the displacement is infinite when its denominator is equal to zero [2]. The evaluation of resonance solutions for flexural waves is

$$
kL = (2n-1)\pi
$$
 (A-18)

Combining the resonance solutions of longitudinal or flexural waves traveling in a finite pile, the resonance solutions are dependent upon their boundary conditions described as Combining the resonance solutions of longitudinal or flexural waves traveling in a finite pile, the resonance solutions are dependent upon their boundary conditions described as

$$
kL = n\pi \text{ for displacement CB} \tag{A-19}
$$

$$
kL = (2n-1)\pi \text{ for stress CB} \tag{A-20}
$$

The resonant wavelength of waves, λ_d , with a displacement CB and resonant wavelength of waves, λ_s , with a stress CB are

$$
\lambda_d = \frac{2L}{n} \tag{A-21}
$$

$$
\lambda_s = \frac{4L}{2n - 1} \tag{A-22}
$$

respectively.

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